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Experimental observation of  
magnetically induced:  
LINEAR DICHROISM ON  
VACUUM.

by

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and myself.

These results have been obtained  
at the PVLAS APPARATUS built  
and operating at the:

LABORATORI NAZIONALI DI LEGNARO  
(L.N.L.);  
LEGNARO, PADOVA ITALY

(2)

← + Breezer



# P.V.L.A.S

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Specifically built in order to obtain significative experimental informations on VACUUM using optical techniques

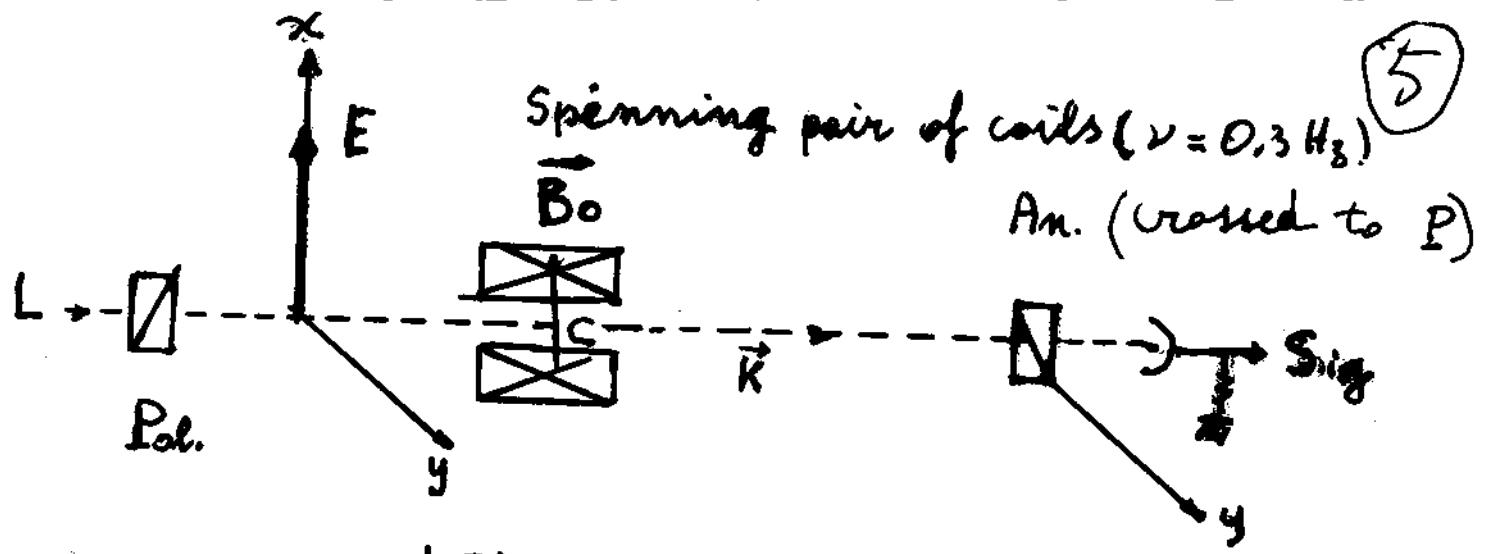
Our full experimental program is to detect and measure the two following properties (eventually) acquired by VACUUM due to the presence of a magnetic field ( $\vec{B}_0$ ):

- 1) Linear Birefringence
- 2) Linear dichroism.

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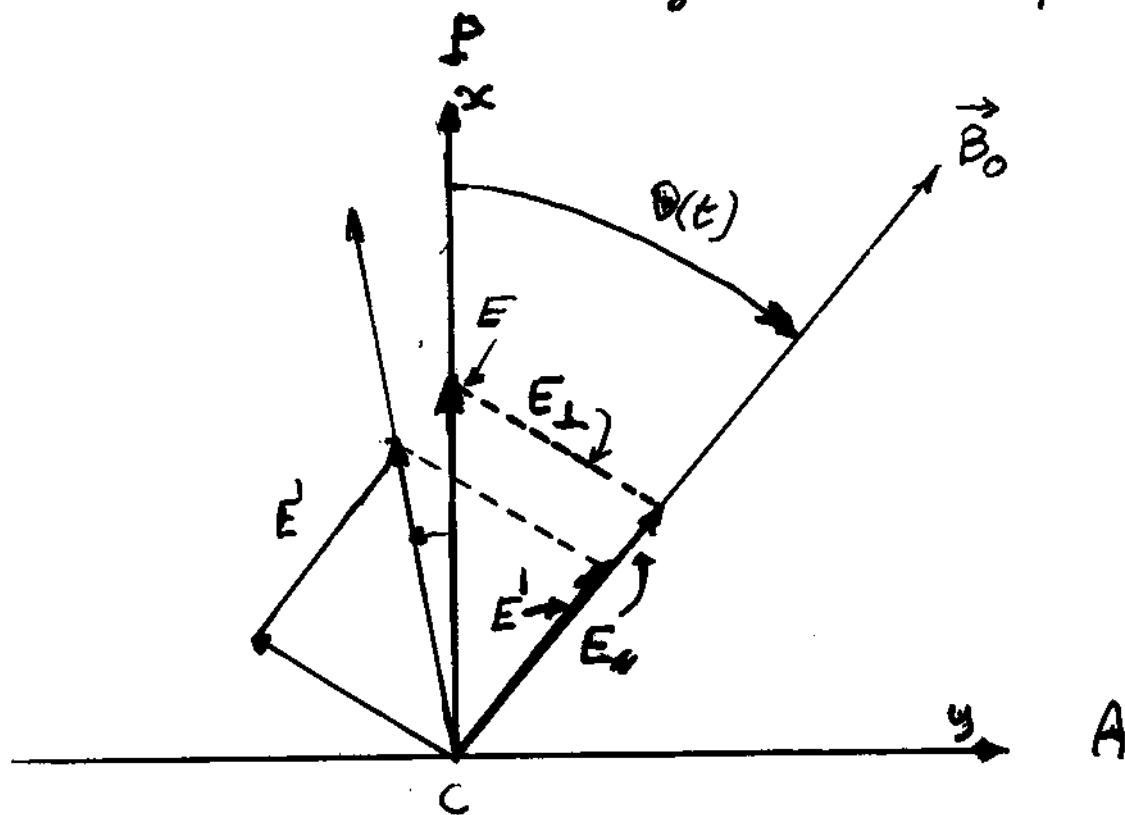
In this communication I report experimental results on  
Linear dichroism induced  
on Vacuum by a static magnetic field ( $\vec{B}_0$ ).

↓  
Selective absorption (by the magnetized vacuum) of polarized light, according to its polarization-orientation with respect to the  $\vec{B}_0$  direction.



if  $|\vec{B}_0| = 0 \rightarrow S_{\text{sig}} = 0$  : beam in vacuum

for  $\vec{B}_0 \neq 0$   $\vec{B}_0$  stays in  $x-y$  plane



Suppose there is a selective absorption of the  $E_{\parallel}$  component. (average value q)

$$E_{\parallel} \rightarrow q E_{\parallel} = E' = E^1 \quad q < 1$$

Initially  
-no absorption-

$$E e^{-i\frac{\epsilon}{2}} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \rightarrow \begin{pmatrix} P \\ A \end{pmatrix}$$

With absorpt.

$$E e^{-i\frac{\epsilon}{2}} \left( \begin{pmatrix} (q-1) \cos^2 \theta(t) + 1 \\ q-1 \sin 2\theta(t) \end{pmatrix} \right)$$

$$\text{Sig } A = \frac{q-1}{2} E \sin 2\theta(t)$$

So max. rot. signal =  $\boxed{\Delta\alpha = \frac{q-1}{2} \text{ rad.}}$

Notice  $\Delta\alpha \neq \Leftrightarrow q \neq 1$

Therefore observing a magnetically induced dichroism in VACUUM means detect a (rotational) signal of a linearly polarized radiation beam that goes through the "magnetized" Vacuum with its propagation vector  $\vec{k}$  orthogonal to  $\vec{B}_0$ .

↓  
OUR experim. results →

With

$$\left\{ \begin{array}{l} L = 110 \text{ cm} \\ \vec{B}_0 = 5.51 \text{ Tesla} \end{array} \right\} \text{ we find}$$

$$\textcircled{1} \quad \boxed{\frac{q-1}{2} = -(3.7 \pm 0.4) 10^{-12} \text{ rad/Pass}}$$

putting  $\bar{N}$  of passes = 52000  
gives:

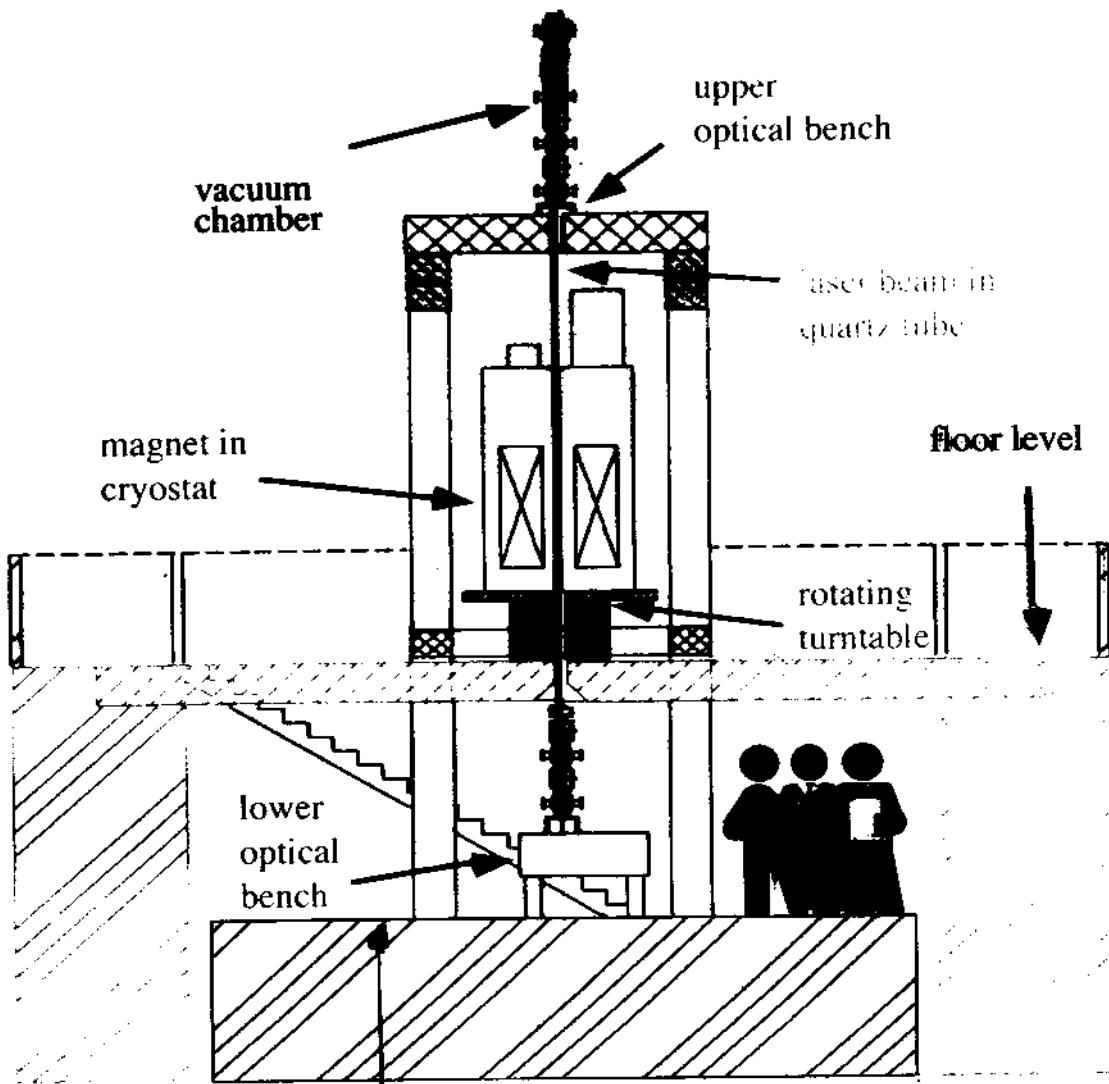
$$\textcircled{1} \quad \boxed{\text{rot. signal} = \Delta \alpha = (2.2 \pm 0.2) 10^7 \text{ rad}}$$

i.e. We observe a change in the amplitude of the polarization component parallel to the  $\vec{B}_0$  field (in fact a decrease).

What happened to the missing part?

- + At this point, before speaking of additional experimental results,
- + compare our to others results,
- + how to proceed to understand what is happening,
- + near future plans,

Let me briefly present the P.V.L.A.S. apparatus.



- The floor of the pit was constructed to be seismically isolated from the rest of the building. It lays on 15 meter pylons sunk in sand.

- The optics is mounted on two granite optical tables: the lower one on the floor of the pit and the upper one on top of a granite structure 8 meters high also resting on the floor of the pit.

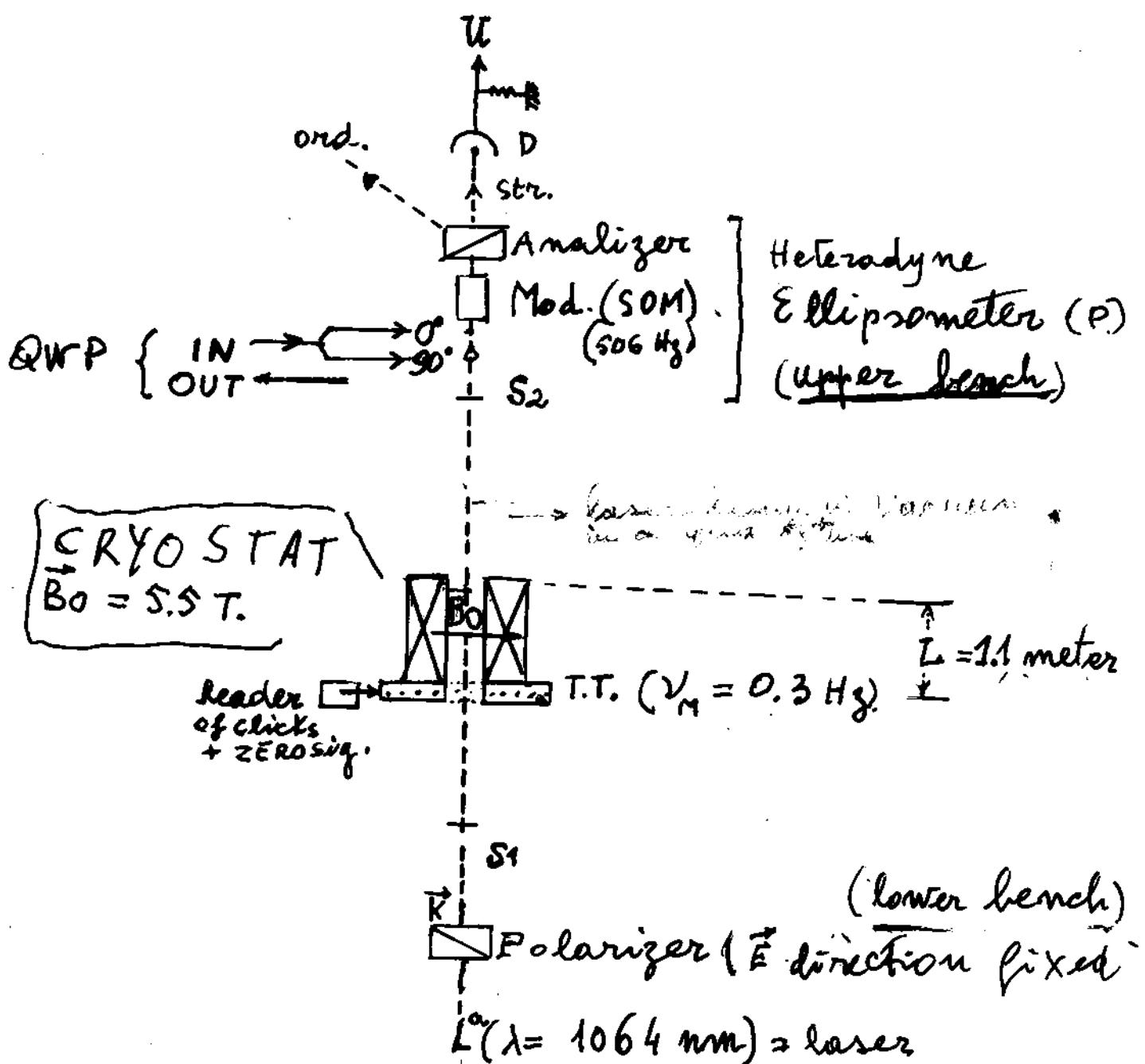
- The cryostat containing the superconducting magnet is fixed on a rotating turntable which in turn is fixed to a reinforced concrete beam laying across the pit but resting on the building floor: when the platform is rotating the field  $B_0$  of the superconducting magnet rotates remaining always on a horizontal plane:  $\nu = 0,3 \text{ Hz}$ .

- As a result the concrete beam and cryostat are mechanically decoupled from all the optical parts.

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# DICHOISM (of VACUUM + $\vec{B}_0$ )

↑ vertical direction :  $\vec{B}_0$  contained in horizontal plane  
while T.T. table rotates.



• [SOM STRESS. OPT. MOD.

• [S<sub>1</sub> S<sub>2</sub> → F. P. cavity locked with laser  
640 cm long

• [T.T. Turning Table  
(with hole center).

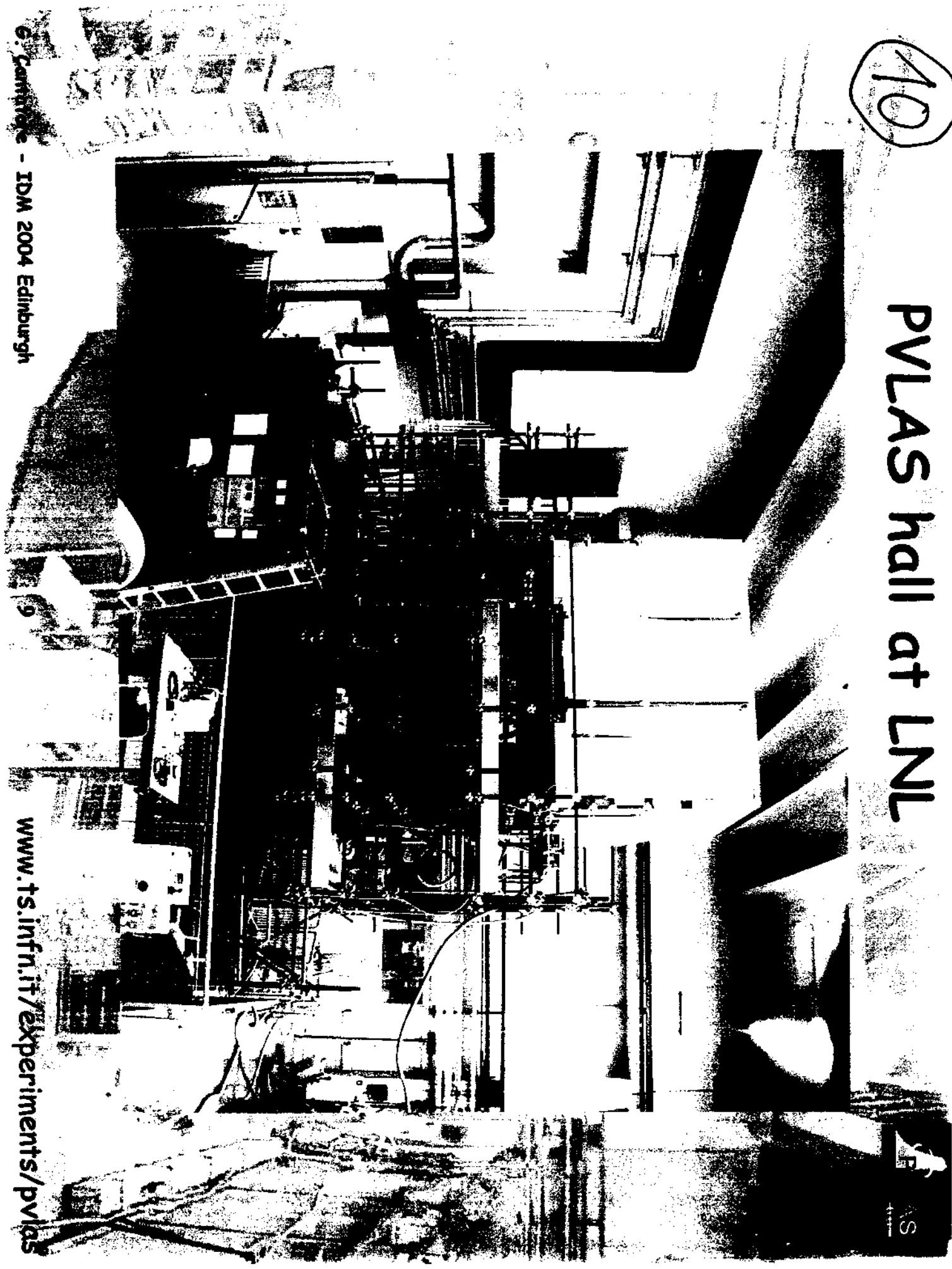
• Laser beam in Vacuum  $< 5.10^{-8} \text{ mb.}$

1 - " Interaction. region. area. 1 mm<sup>2</sup>.

fig. 10

# PVLAS hall at LNL

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## MAIN PARAMETERS OF THE PVLAS APPARATUS

- MAGNET

DIPOLE, 6T, Temp. 4.2 K,  $1 \text{ m FIELD ZONE} = L_m$   
 $B_0 \rightarrow \text{HORIZONTAL}$

VERY IMPORTANT

- CRYOSTAT

ROTATION FREQ. (TYPICAL)  $\sim 300 \text{ mHz}$ , SLIDING CONTACTS, WARM BORE

- LASER

1064 nm, 100 mW, FREQUENCY-LOCKED TO THE FABRY-PEROT CAVITY

- FABRY-PEROT OPTICAL CAVITY

$6.4 \text{ m LENGTH}$ , FINESSE  $\sim 10^5$ , OPTICAL PATH IN THE INTERACTION REGION  $\sim 60 \text{ Km}$

$S_1 S_2$

- HETERODYNE ELLIPSOMETER

S.O.M. = ELLIPTICITY MODULATOR, QUARTER-WAVE PLATE (ON-OFF), HIGH EXTINCTION ( $\sim 10^{-7}$ )

- SOM provides ellipticity carrier for Heterodyne
- DETECTION CHAIN.

PHOTODIODE WITH LOW-NOISE AMPLIFIER

- DAQ

DEMODULATED AT LOW FREQUENCY AND PHASE-LOCKED TO THE MAGNETIC FIELD INSTANTANEOUS DIRECTION

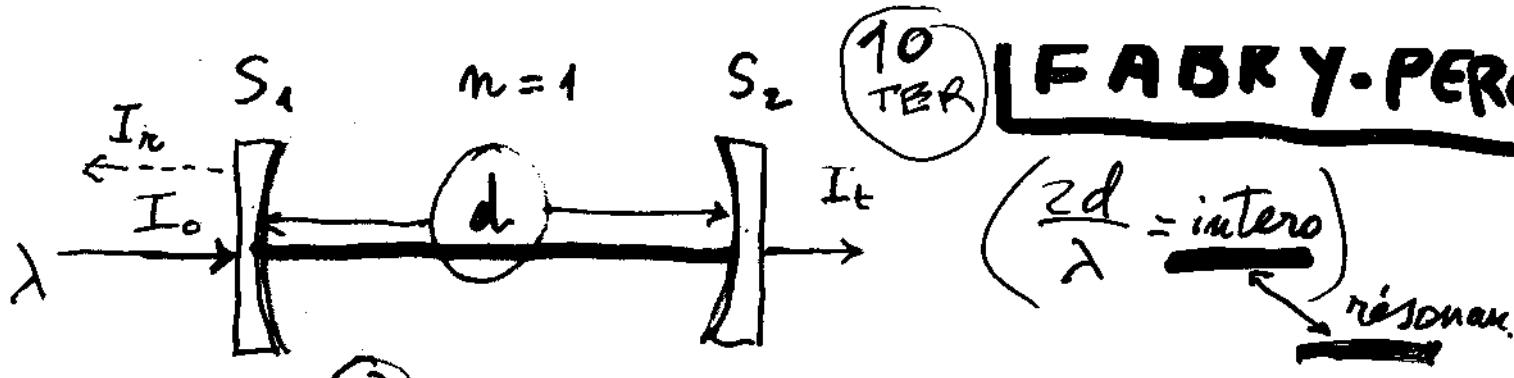
HIGH SAMPLING FREQUENCY (8.2 kHz)

DIRECT ACQUISITION

- DICROISM CONFIGURATION QWP in  $\{ \begin{matrix} 0^\circ \\ 90^\circ \end{matrix} \}$

- Ellipticity " "

QWP out



$(T, R, \eta, R_c)$

$$\left(\frac{I_t}{I_o}\right)_t = \eta_t^o \frac{1}{1 + \left(\frac{2F}{\pi}\right) \sin^2 \delta} \cdot \left(\frac{I_r}{I_o}\right)$$

$$\lambda = 1060 \text{ nm} \quad \nu = 3 \times 10^{14} \text{ Hz}$$

$$\delta = 2\pi \frac{2d}{\lambda}$$

$$\eta_t^o = \left(\frac{TF}{\pi}\right)^2$$

$$F = \frac{2\pi c}{2d} \quad (C)$$

fineness  $\dots \rightarrow \approx \text{reflectivity}$   
 $(\geq 10^5)$   
 $= R \Rightarrow \frac{2}{\pi} F -$

$$\eta_t^o = \left(\frac{TF}{\pi}\right)^2$$

ratio  $I_t/I_o$  at reson.  
 $> 10\%$

$$(cav. \text{ airy}) \quad \Delta V_c = \frac{1}{2\pi} \quad (C)$$

cavity linewidth  
 $(\sim 300 \text{ Hz} \pm 400 \text{ Hz})$

$$Q = \frac{\text{V. height}}{\Delta V_c}$$

quality factor  
 $(\sim 10^{12})$

$$\nu_{FSR} = \frac{c}{2d} \quad (C)$$

free spectral range  
 $(25 \text{ MHz})$

$$\text{Th. } \left\{ \begin{array}{l} F = \frac{\pi \sqrt{R}}{1-R} \\ \eta_t^o = \left(\frac{TF}{\pi}\right)^2 \end{array} \right.$$

$$\begin{array}{l} \sim 10^{-5} \\ \boxed{T+R+q=1} \end{array} \rightarrow \text{perdite}$$

$$\eta_t^o = \left(\frac{TF}{\pi}\right)^2$$

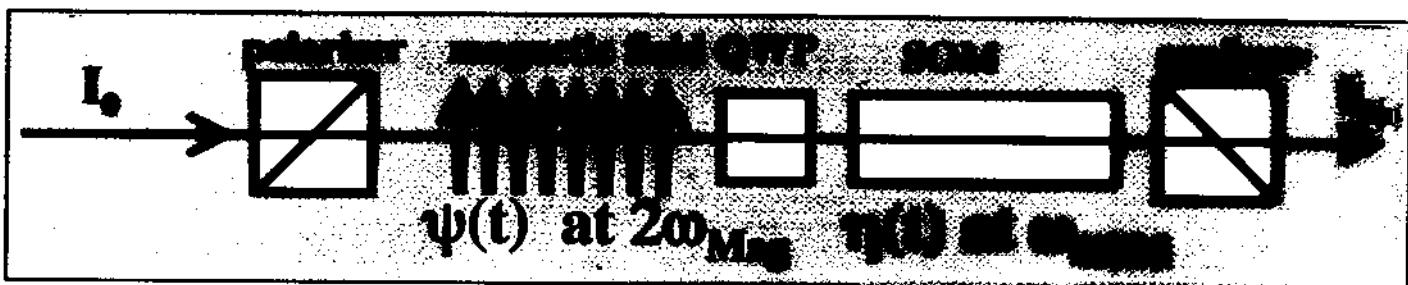
$$N_{\text{refl.}} = \frac{2}{\pi} F \approx \frac{2}{\pi}$$

(1)

For Dichroism compensation

# Measurement principle

- Static measurement is excluded.
- Modulate the field and add a carrier signal with a modulator at  $\omega_{SOM}$
- Rotating the field at  $\Omega$  produces an ellipticity at  $2\Omega$  (B<sup>2</sup> dependence) and  $\Delta\omega$  at  $2\Omega$



Ideally

$$\begin{aligned} I_{tr} &= I_0 \left[ \sigma^2 + (\Psi(t) + \eta(t))^2 \right] - \\ &= I_0 \left[ \sigma^2 + (\Psi(t)^2 + \eta(t)^2 + 2\Psi(t)\eta(t)) \right] \end{aligned}$$

Main frequency components:

- $2\omega_{SOM}$  from  $\eta(t)^2$
- $\omega_{SOM} \pm 2\Omega_{Mag}$  from  $2\Psi(t)\eta(t)$

$$2\sin a \sin b = [\cos(a-b) - \cos(a+b)]$$

$$I_{Tr} = I_0 \left[ \sigma^2 + (\Psi(t) + \eta(t) + \alpha_s(t))^2 \right]$$

$$= I_0 \left[ \sigma^2 + (\eta(t)^2 + 2\Psi(t)\eta(t) + 2\alpha_s(t)\eta(t) + \dots) \right]$$

**Normalization**

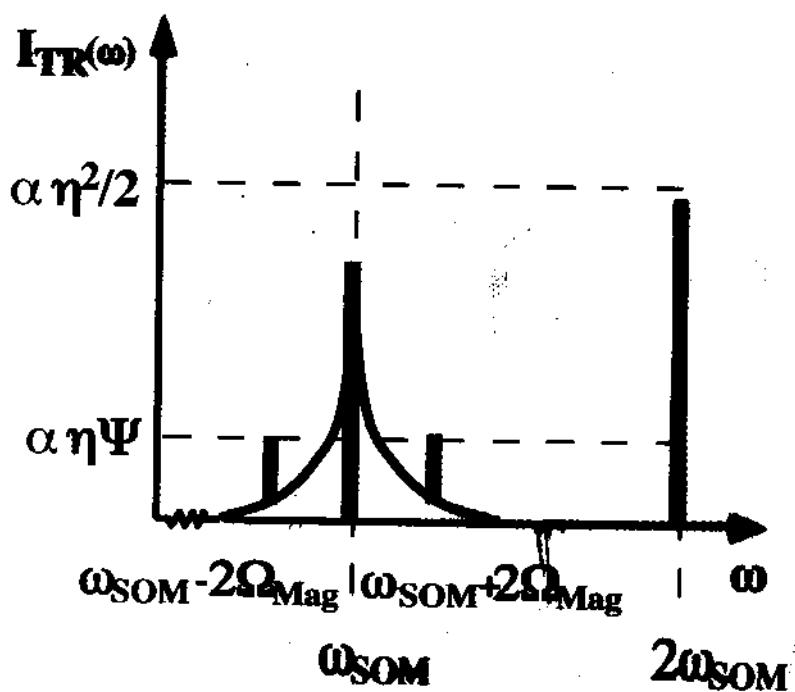
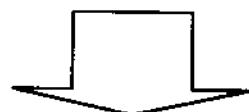
**Desired signal**

**Birefringence noise**

$$\Psi(t) = \Psi_0 \sin(2\Omega_{Mag} t); \quad \eta(t) = \eta_0 \sin(\omega_{SOM} t + \varphi)$$

$$2\Psi(t)\eta(t) = \Psi_0 \eta_0 \{ \cos((\omega_{SOM} - 2\Omega_{Mag})t + \varphi) - \cos((\omega_{SOM} + 2\Omega_{Mag})t + \varphi) \}$$

A small, time-varying signal can be extracted from a large noise background with the heterodyne technique.



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In order to verify that everything is working as due (and under control) we measured the binifrequency of same gasses ( $N_2$ ,  $Ne$ ) and found in agreement to where to appear in a polar plot as far as regards

their PHASES

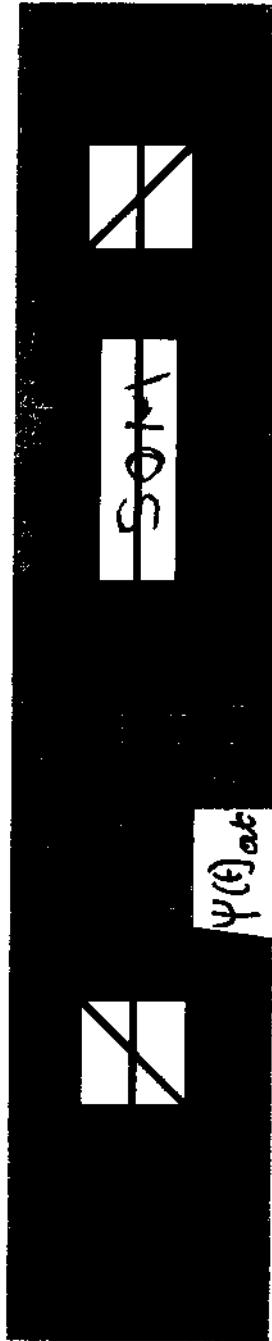
Notice  $N_2$  appont to  $Ne$



→ ELLIPTICITY CONFIGURATION

## Measurement principle

- Static measurement is excluded.
- Modulate the field and add a carrier signal at  $\omega_{\text{SOM}}$
- Rotating the field at  $\Omega$  produces an ellipticity at  $2\Omega$



Ideally,  $\psi(t) = \psi_0 \sin(\omega t)$   
 $45^\circ \quad 90^\circ$

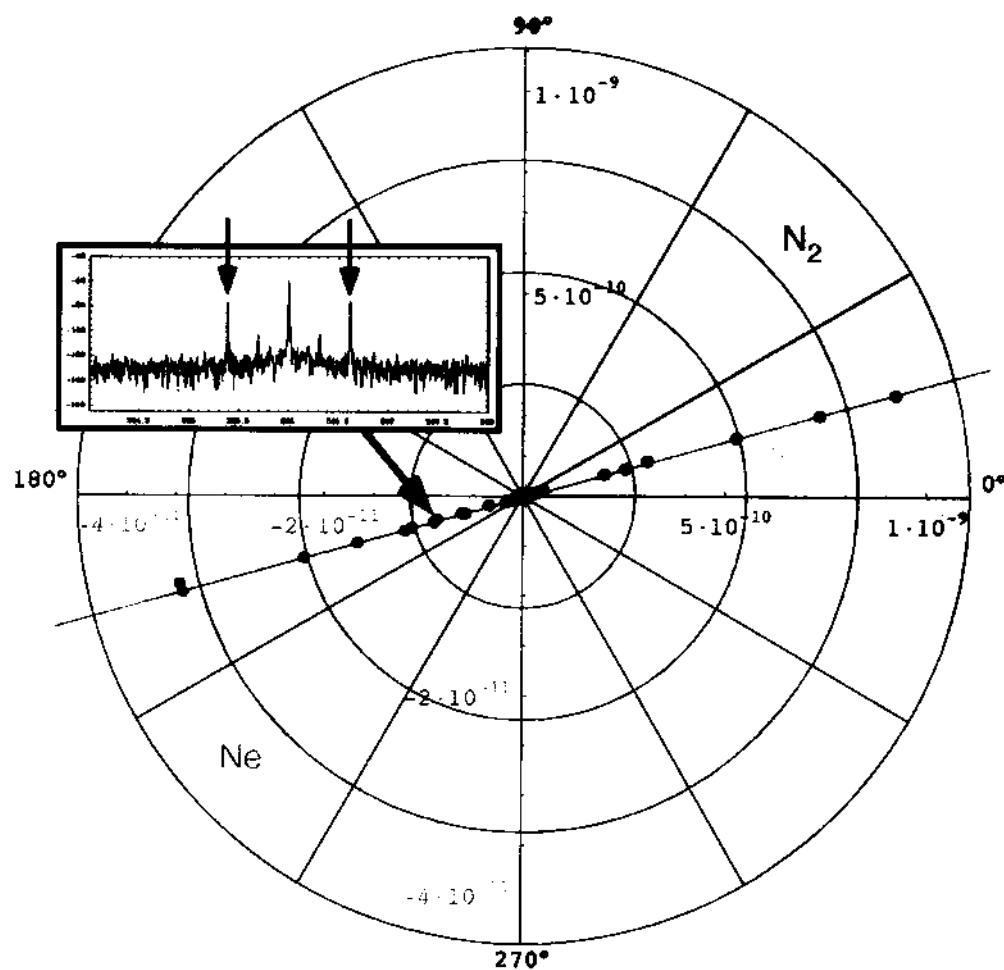
$$I_{Tr} = I_0 [\sigma^2 + (\Psi(t) + \eta(t))^2] = I_0 [\sigma^2 + (\Psi(t)^2 + \eta(t)^2 + 2\Psi(t)\eta(t))]$$

Main frequency components at  $\omega_{\text{SOM}} \pm 2\Omega$  Mag and  $2\omega_{\text{SOM}}$

SOM = Stress Optic Modulation

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- PHYSICAL INFORMATION IS EXTRACTED  
BY ANALYSING AMPLITUDE AND PHASE  
OF THE FREQUENCY COMPONENTS OF  
THE LIGHT TRANSMITTED THROUGH  
THE ELIPSOMETER

# Sensitivity

**Shot noise sensitivity for PVLAS is**

$$\left\{ \sqrt{\frac{4e}{I_0 q}} = 2 \cdot 10^{-8} / \sqrt{\text{Hz}} \quad \approx 150 \text{ mV} \right.$$

*+ altri (parali.)*

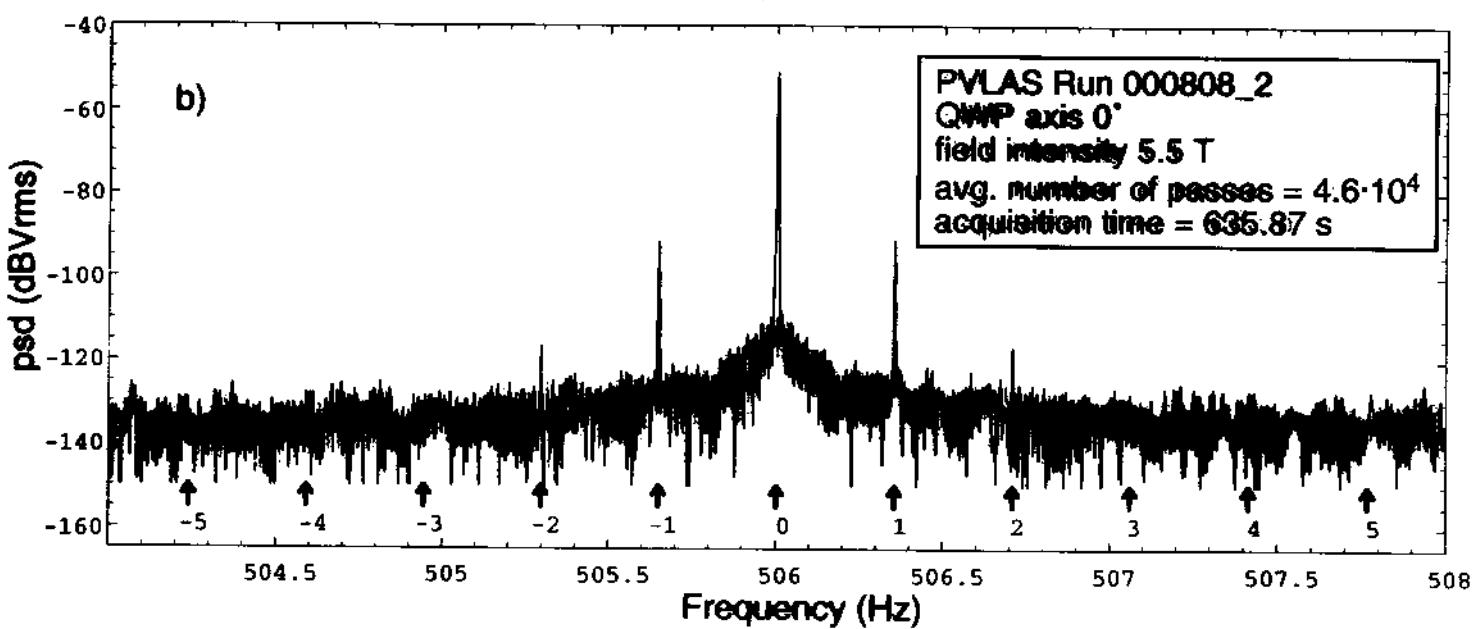
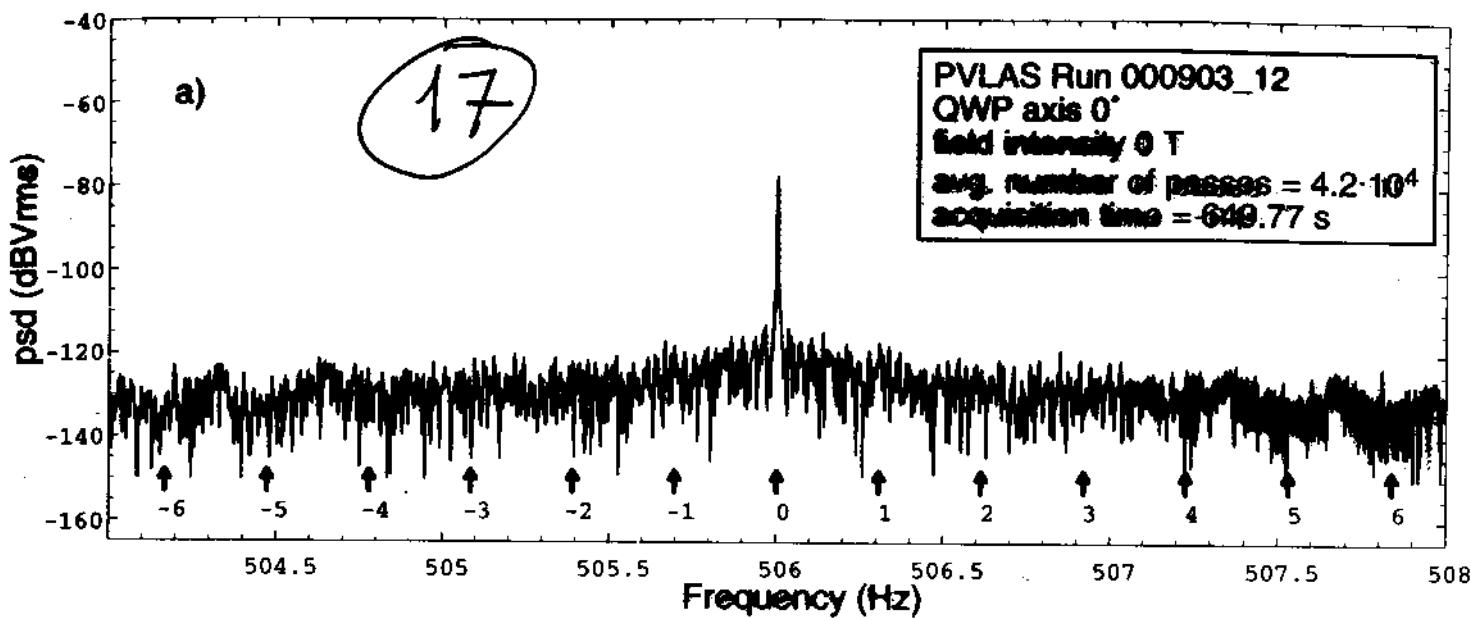
**Present sensitivity for PVLAS is**

$$\approx 2 \cdot 10^{-7} / \sqrt{\text{Hz}}$$

*(continuo snello 3 dB  
dovuto all'effetto di rumore  
lineare)*

**Dichroism signal from PVLAS is**

$$= (2.2 \pm 0.3) \cdot 10^{-7} \text{ rad}$$



Signals between 506 & define

(18)

$2.5 \cdot 10^7$

$2 \cdot 10^7$

a)

$1.5 \cdot 10^7$

$1 \cdot 10^7$

$5 \cdot 10^6$

506.4

506.6

506.8

507

507.2

507.4

Frequency (Hz)

PVLAS Run 659\_1

QWP axis 0°

field intensity 0 T

avg. number of passes =  $4 \cdot 10^4$

acquisition time = 715.50 s

rotation (rad)

$2.5 \cdot 10^7$

$2 \cdot 10^7$

b)

$1.5 \cdot 10^7$

$1 \cdot 10^7$

$5 \cdot 10^6$

506.4

506.6

506.8

507

507.2

507.4

Frequency (Hz)

PVLAS Run 657\_2

QWP axis 90°

field intensity 5.5 T

avg. number of passes =  $4 \cdot 10^4$

acquisition time = 654.38 s

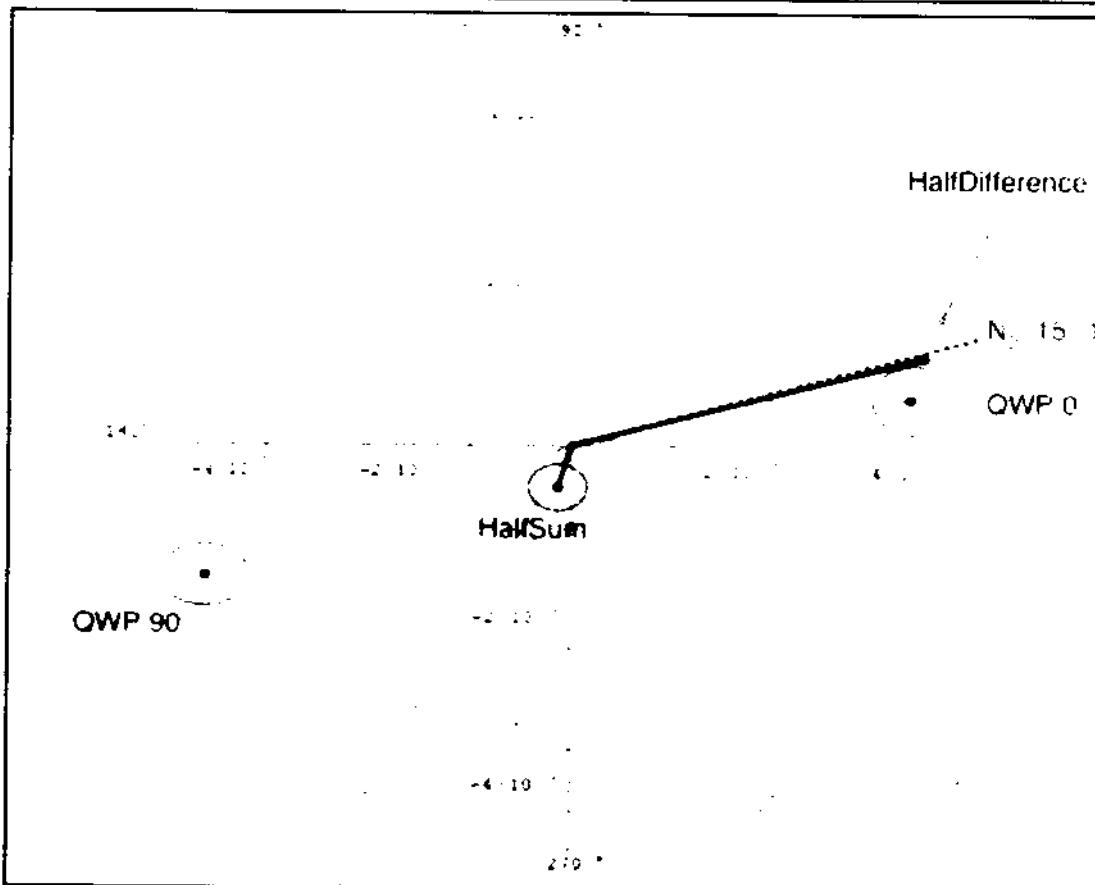
(18)

## Results for the measured dichroism in vacuum

- QWP 90°: Result of weighted average with the quarter wave plate at 90°
- QWP 0°: Result of weighted average with the quarter wave plate at 0°
- ..... N<sub>2</sub>: Physical axis defined by measuring the Cotton-Mouton effect in Nitrogen

— HalfDifference:  $\frac{\text{QWP}90^\circ - \text{QWP}0^\circ}{2} = \text{Dichroism}$

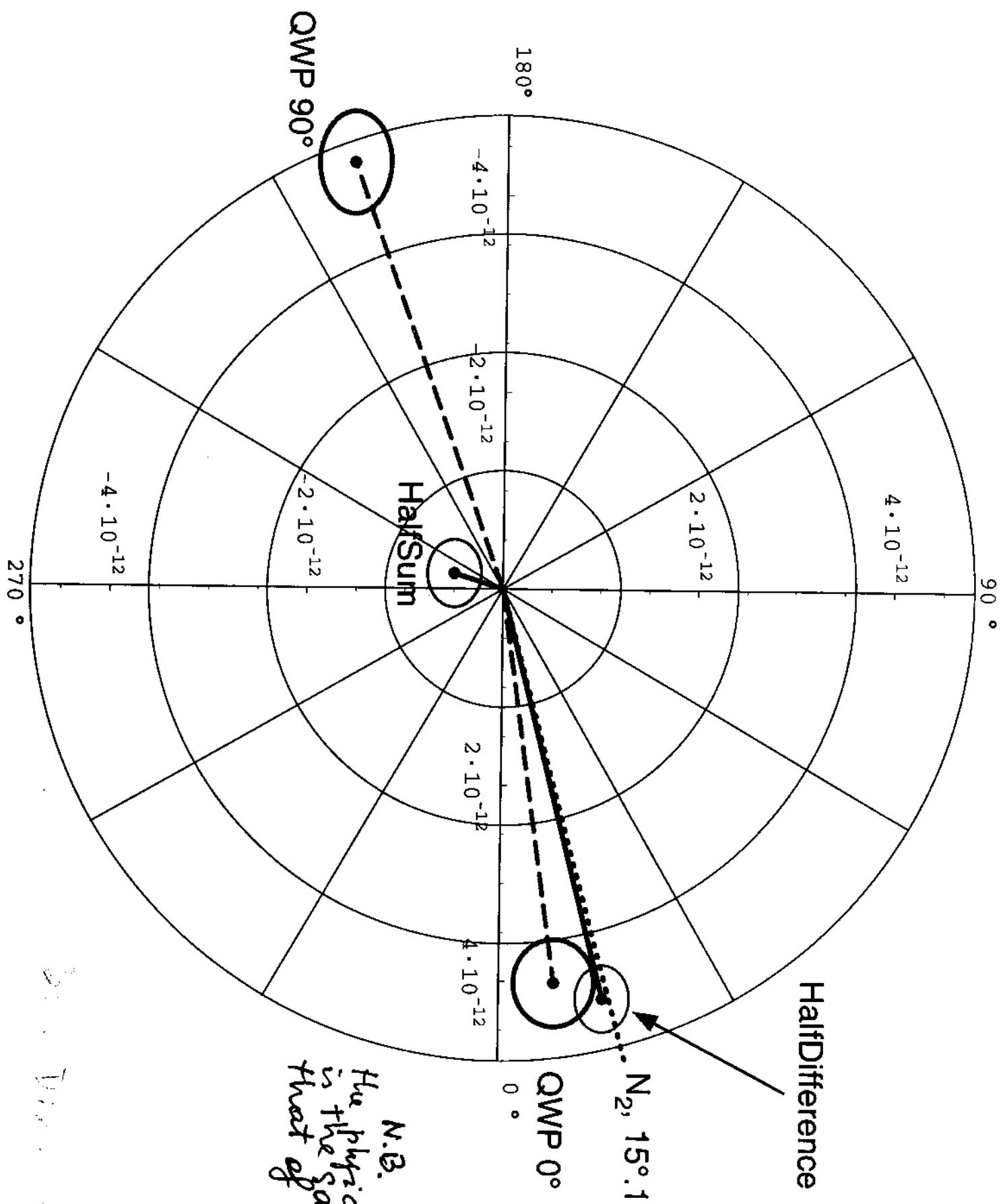
— HalfSum:  $\frac{\text{QWP}90^\circ + \text{QWP}0^\circ}{2} = \text{Spurious signal}$



$$\text{Dichroism signal} = (2.2 \pm 0.3) \cdot 10^{-7} \text{ rad} = \Delta \alpha$$

The phase of the vacuum dichroism equals the phase of the Nitrogen Cotton-Mouton effect: both in turn agree with the expect axis of asymmetry

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# Final Result

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$$\Delta \alpha = -\frac{q-1}{2} = (2.2 \pm 0.3) 10^{-7} \text{ rad}$$

with the conditions

$$l_m = 110 \text{ cm}$$

$$N = 52000 = \frac{e}{n} F$$

$$|B_0| = 5.51 \text{ Tesla}$$

$$F = \frac{\pi c \tau}{d_{Mi}} : \tau \text{ measured}$$

$$420 \leq \tau \leq 620 \mu\text{s}$$

[accordo accordinga di accordo con B.N.L. ?]

QUESTION what is the origin of  $q \neq 1$  ?

- 1) Possible unexpected systematic mistake
- 2) if real: is it an absorption? a mixing? (of photons).

We have planned a "near-future" shut down in order to rebuilt the apparatus correcting some malfunctions that we have learned.

Among others:

- a) Sincere general insulations
  - b) correct, via optical feed back, the residual beam movements: improve the housing of the optical pieces
  - c)  $\lambda \rightarrow \lambda' = \frac{\lambda}{2}$  (532 nm) "Green" (one can)
  - d) improve the phonics-shielding
- 3) May be we seeing a strong QED-QCD interference - ending with photon splitting -
  - 4) DECIDED The most important thing now is determine  $[m_e]$  to see if our apparatus is O.K.

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Hypothesis  
~~absorption~~

Mixing (particle of)  
(mass  $m_{ac^2}$ )  
↓

We have an  
expression to  
which compare  
the obtained  
results

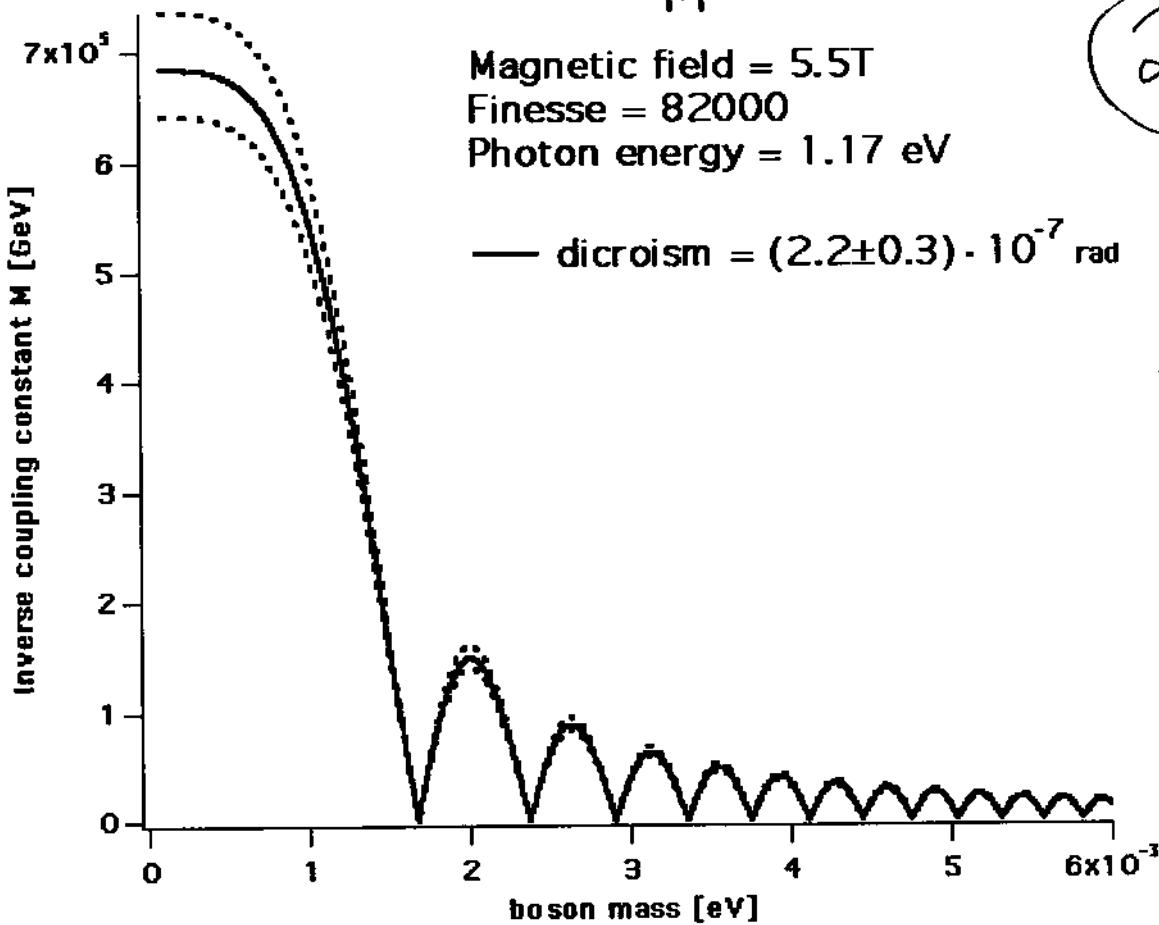
However test of  
coherence

$$M_b < \frac{27 w}{L_M}$$

$$\downarrow \\ \sim 1.3 \cdot 10^{-3} \text{ ev}$$

$\downarrow$   
method to find  $m_{ac^2}$

Hypothesis: the observed rotation due to a mixing of the photons (of the beam) with a light boson (mass =  $m_\gamma c^2$ ) neutral and with a coupling to two photons  $\frac{1}{M} \text{ GeV}^{-1}$



Natural Heaviside-Lorentz units

$$\Delta \alpha = \epsilon = -\sin(2\alpha) \left( \frac{BL}{4M} \right)^2 N \left( \frac{\sin x}{x} \right)^2 \quad \leftarrow \begin{matrix} \text{rotation} \\ \text{Signal} \end{matrix}$$

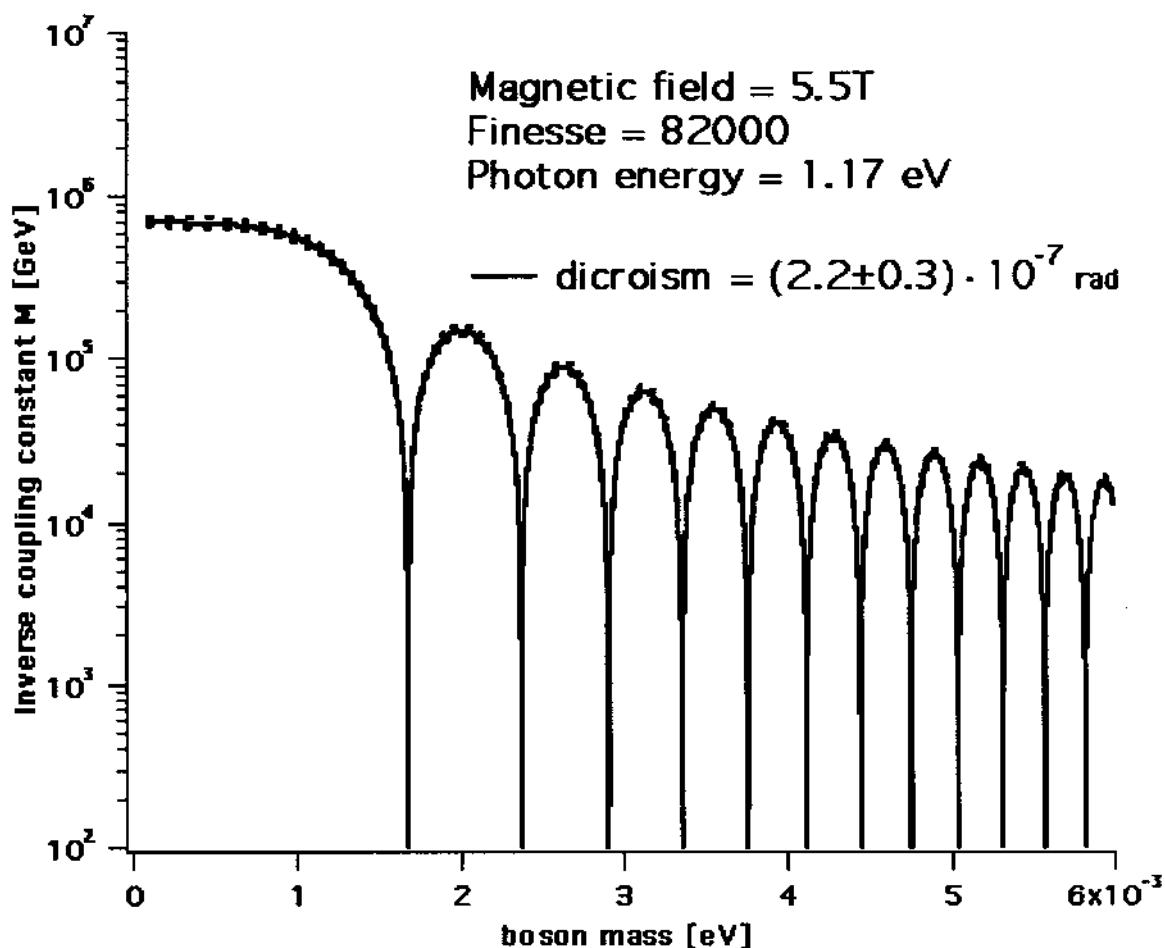
$$x = \frac{L}{2} \left[ \frac{k_m^2}{2k} + (n-1)k \right]$$

$$k = \frac{2\pi}{\lambda}; \quad N = \frac{2F}{\pi}; \quad k_m = \frac{mc}{\hbar}$$

if  
due to  
a boson  
mixing with  
 $\gamma$

$m_\gamma c^2$  = mass of the boson.

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Natural Heaviside-Lorentz units

$$\bar{N} = \frac{3}{\pi} \text{ Finesse} \approx 52000$$

Expecting to obtain scaling data  
in different  $\sim 2000$ .

① Adapt expression for  $\Delta\alpha^0$  coming from the hypothesis of "existing" e light BOSONS (of mass  $M_e c^2$ ), neutral, coupled to two photons with a C.C.  $\frac{1}{M} \text{ GeV}^{-1}$

$$\Delta\alpha^0 = \left( \frac{BL}{4M} \right)^2 \left( \frac{\sin x}{x} \right)^2 \cdot N \quad \text{(A)}$$

$$x = \frac{L}{2} \left[ \frac{k_m^2}{2K} + (n-1)K \right]$$

$$k = \frac{2\pi}{\lambda}$$

$$N = \frac{2}{\pi} F$$

$$k_m = \frac{mc^2}{\pi c}$$

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B15

$n$  = index of refraction (1 atm.) of neon

$$n-1 = 6.7 \cdot 10^{-5}$$

② take many data with different pressure of neon

③ do fits of expression (A) after subtracted the signal due to the Cotton-Mouton effect of neon (component proportional to pressure)  
thus to deduce

$M_e c^2 = \text{mass of the "hypothetical boson"}$

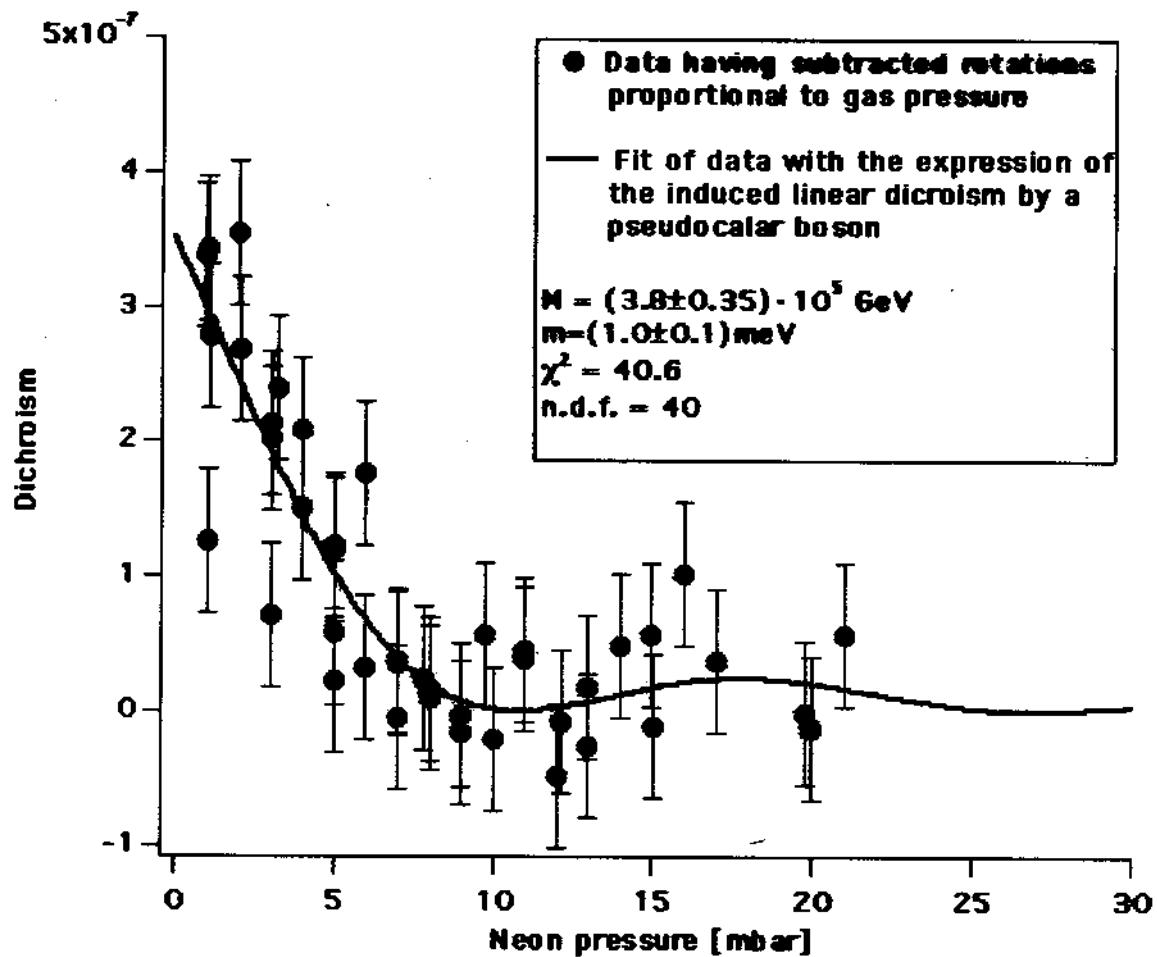
↳ these produce the test of coherence for our apparatus (see preceding page)

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B15

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## Dichroism Vs. Neon pressure

$Q \neq P \quad \theta$



$$M = (3.8 \pm 0.35) \cdot 10^5 \text{ GeV}$$

$$m = (1.0 \pm 0.1) \text{ meV} \rightarrow [(1.0 \pm 0.1) < 1.3] \cdot 10^3 \text{ el} \\ - \Theta. K. -$$

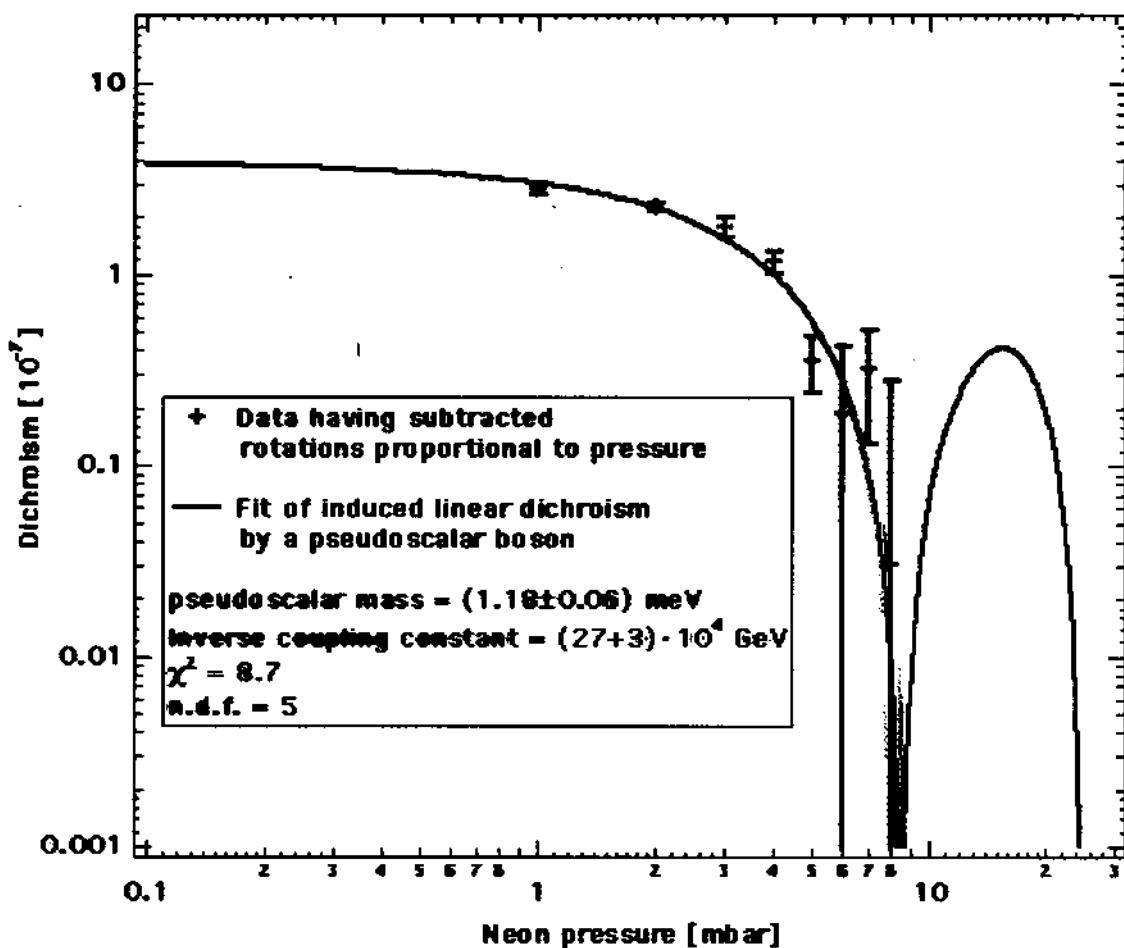
- $P_{\text{Neon}} = 0 \Rightarrow \text{Dichroism} = (2.9 \pm 0.4) \cdot 10^{-7} \text{ rad}$
- Vacuum dichroism =  $(2.2 \pm 0.3) \cdot 10^{-7} \text{ rad}$

all referred to number of passes  $\approx 52000$ .

more surprising thing!

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## Dichroism Vs. Neon pressure



QWP 0°

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Universita' degli studi di Ferrara  
Dipartimento di Fisica



The cryostat showing the quench protection system.

$L_M$  small to preserve the relative phase between the light boson and the photon's field.

$$m_b < M_0^2 = \frac{2\pi c}{L_M}$$

$$M_0^2 \approx (1.3 \cdot 10^{-3} \text{ eV})^2$$

Tucson, Arizona 2003

inequality to be verified.



- general  
discrepancy  
resistance

$L_{MAGNET} = 11$

⊗ ⊗ ⊗

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$\nu_{\alpha}$ [eV]	$H$ [GeV]	$g_{eff} = \frac{1}{M} [GeV^{-1}]$
$(1.0 \pm 0.1) \times 10^{-3}$	$(3.8 \pm 0.4) \times 10^5$	$(2.6 \pm 0.5) \times 10^{-6}$
PVLAS Ne DATA		
$(5 \pm 1) \times 10^{-5}$	$(2.0 \pm 0.4) \times 10^{-6}$	
PVLAS VACUUM ROTATION		
$\left[ (1.0 \pm 0.1) \times 10^{-3} \right]$	$\geq (4 \times 10^{-5})$	$\leq 2.5 \times 10^{-6}$
BFR T [PRD 47, 9 (1993)]		
ELLIPTICITY DATA ONLY		
- Compatible with a photon-hadron mixing mechanism $M_{\pi} c^2 \approx (1.0 \pm 0.1) MeV ; M^{\pi} = (3.8 \pm 0.35) 10^5 GeV$		

# International Workshop on “Neutrino Telescopes”

① We observe in vacuum

a polarization rotation signal

$$\Delta\phi \neq 0$$

② a possible interpretation  
is to assume the existence  
of a very light, neutral,  
pseudoscalar boson coupled  
to two photons. TEST coherence O.K.

③ With these Hypotheses  
we obtain

$$m_\phi = 1 \pm 0.1 \text{ meV}$$

$$g_{\phi\gamma} = (2 \div 3) 10^{-6} \text{ GeV}^{-1}$$

④ We plan to confirm these  
results in the near future  
by using energy doubled  
laser photons.

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